

Non-linear damping of kink waves

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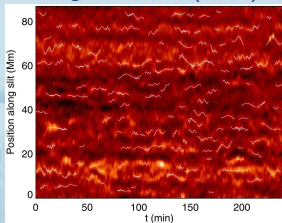
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Standing and propagating kink waves in solar corona.
Observations of non-linear damping.

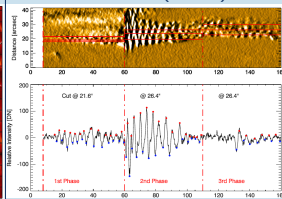
Propagating
transverse waves in
the coronal loops
and plumes

Thurgood et al. (2014)

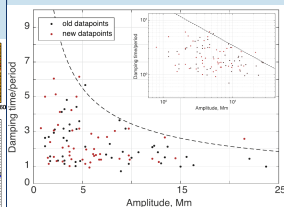


Standing transverse
waves (after flares &
decayless)

Nistico et al. (2013)



Damping of standing
transverse waves is
amplitude dependent



Nechaeva et al. (2019):

$$\tau/P \sim A^{-0.7}$$



Non-linear damping through KHI and uniturbulence

Previous simulations of standing kink waves show KHI

Propagating kink waves have uniturbulence

Cascade to small scales, loss of energy from mode \rightarrow damping.

Elsässer variables:

$$\vec{z}^{\pm} = \vec{v} \pm \frac{\vec{b}}{\sqrt{\mu\rho}}$$

Governing equations:

$$\frac{\partial \vec{z}^{\pm}}{\partial t} \mp \vec{v}_A \cdot \nabla \vec{z}^{\pm} = -\vec{z}^{\mp} \cdot \nabla \vec{z}^{\pm}$$

Uniturbulence:

$$(\omega + \omega_A) \vec{z}_{\perp}^{+} = (\omega - \omega_A) \vec{z}_{\perp}^{-}$$

Edwin & Roberts (1983):

Standing kink waves

$$P'(r, \varphi, z, t) = \mathcal{R}(r) \cos \varphi \cos(k_z z) \cos(\omega t)$$

Propagating kink waves

$$P'(r, \varphi, z, t) = \mathcal{R}(r) \cos \varphi \cos(k_z z - \omega t)$$

with

$$\mathcal{R}(r) = \begin{cases} A \frac{J_1(\kappa_i r)}{J_1(\kappa_i R)} & \text{for } r \leq R, \\ A \frac{K_1(\kappa_e r)}{K_1(\kappa_e R)} & \text{for } r > R, \end{cases}$$

	MHD turbulence		Uniturbulence	
	Upward	Downward	Upward	Downward
z^{-}	✓		✓	
z^{+}		✓	✓	
	counterpropagating $\omega = \pm \omega_A$		co-propagating $\omega \neq \pm \omega_A$	

Damping of propagating waves

RMS average z^2 and $\epsilon = \vec{z} \cdot \nabla w$ over period, wavelength, cross-section

Energy density

$$\langle w \rangle = \pi R^2 \frac{\rho_i + \rho_e}{2} V^2$$

$$\langle \epsilon \rangle = V^3 \frac{\sqrt{5\pi} R}{10} \frac{\rho_e}{\omega^3} |\omega(\omega^2 - \omega_{Ae}^2)|$$

Damping time:

$$\tau = \sqrt{5\pi} \frac{R}{V} \frac{2(\zeta+1)}{|\zeta-1|}$$

$$= \sqrt{5\pi} \frac{P}{2\pi a} \frac{2(\zeta+1)}{|\zeta-1|}$$

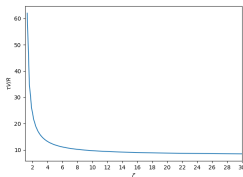
with density contrast

$\zeta = \rho_i / \rho_e$, velocity

amplitude V and

maximal displacement

$\eta = aR$



Examples:

- Pant et al.

(2019):

$V = 22\text{km/s}$,

$R = 250\text{km}$,

$\zeta = 3$

$\rightarrow \tau \sim 180\text{s}$

- plumes with

$R = 1\text{Mm}$,

$V = 4\text{km/s}$,

$\zeta = 3$:

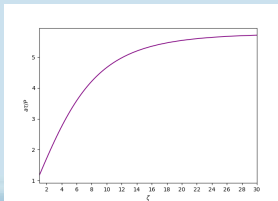
$\rightarrow \tau \sim 3960\text{s}$

Damping of standing waves

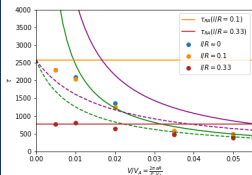
Damping time:

$$\tau = 20\sqrt{\pi} \frac{R}{V} \frac{1+\zeta}{\sqrt{\zeta^2 - 2\zeta + 97}}$$

$$= 20\sqrt{\pi} \frac{P}{2\pi a} \frac{1+\zeta}{\sqrt{\zeta^2 - 2\zeta + 97}}$$

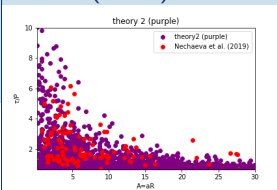


Blue, yellow and red dots: simulations of Magyar & VD (2016)



Full line: non-linear damping
Horizontal: resonant absorption
Dashed: harmonic average

Red dots: Nechaeva et al. (2019)



Purple: 5000 draws from $\zeta \in U[1, 9.5]$,
 $I/R \in U[0, 2]$,
 $A \in U[0.2, 30] \text{Mm}$,
 $R \in U[0.5, 5] \text{Mm}$

Theoretically:

A^{-1} power law